INDIRECT ABELIAN THEOREMS AND A LINEAR VOLTERRA EQUATION

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1. Introduction and summary. We study asymptotic behavior of solutions of

(1.1)
$$x'(t) = k - \int_0^t [a(t-\tau) + c] x(\tau) d\tau, \qquad x(0) = x_0, \qquad \left(' = \frac{d}{dt} \right)$$

where k and x_0 are real, $c \ge 0$, and a(t) satisfies

(H1) $a(t) \in C(0, \infty) \cap L_1(0, 1)$. a(t) is nonnegative and nonincreasing, $\lim_{t\to\infty} a(t) = 0$, and $0 < a(0+) \le \infty$;

(H2) a(t) is convex downward; i.e., for $0 < \varepsilon < 1$ and $0 < t_1 < t_3 < \infty$, $\varepsilon a(t_1) + (1 - \varepsilon)a(t_3) \ge a(t_2)$, where $t_2 = \varepsilon t_1 + (1 - \varepsilon)t_3$.

By a familiar theorem on Volterra equations, (1.1) has a unique solution in $C^1[0, \infty)$. We define u(t) as the solution of (1.1) with k=0, $x_0=1$, and we let $w(t) = \int_0^t u(\tau) d\tau$. It is easily checked that the solution of (1.1) is given by $x_0 u(t) + kw(t)$.

Treating (1.1) as a special case of a nonlinear equation, Levin proved in [5] that if $a(t) \in C[0, \infty)$, $a(t) \not\equiv a(0)$, and $(-1)^k a^{(k)}(t) \ge 0$ for $0 < t < \infty$, k = 0, 1, 2, 3, then

$$\lim_{t\to\infty}u(t)=0$$

and

(1.3) (i) If $c+a(t) \in L_1(0, \infty)$ (in particular, c=0), then

$$\lim_{t\to\infty}w(t)=\left(\int_0^\infty a(t)\,dt\right)^{-1}$$

(ii) If c > 0, then $\lim_{t\to\infty} w(t) = 0$.

Levin also conjectured that

(1.4) If
$$c = 0$$
, $a(t) \notin L_1(0, \infty)$, then $\lim_{t\to\infty} w(t) = 0$.

The present theorem shows that (H1) and (H2) together are nearly sufficient for (1.2), (1.3), and (1.4); in particular, Levin's conjecture is proved. The theorem also exhibits a class of kernels satisfying (H1) and (H2) but for which a different asymptotic behavior from (1.2), (1.3), and (1.4) can be established.

More specifically, note that if a(t) is given by

(a)
$$a(t) = \sum_{k=1}^{\infty} \delta_k \left(1 - \frac{\min\{t, kt_0\}}{kt_0} \right), \quad t_0 = \frac{2\pi}{\tau_0} > 0$$

(1.5) (b)
$$\delta_k \ge 0, \qquad 0 < \delta \equiv \sum_{k=1}^{\infty} \delta_k = a(0) < \infty$$

(c)
$$\Omega \equiv \{k \mid \delta_k > 0\}$$
 has no common divisor > 1,

then $a(t) = \sum_{j=k}^{\infty} \delta_j - t \sum_{j=k}^{\infty} (\delta_j | jt_0)$, $(k-1)t_0 \le t \le kt_0$. It follows that a(t) is continuous, and on each interval $(k-1)t_0 \le t \le kt_0$ it is linear with slope $-\sum_{k=j}^{\infty} (\delta_j | jt_0)$. Then a(t) satisfies (H1) and (H2). When (1.5) holds, we may also have

(1.6)
$$\omega \equiv \sqrt{(\delta + c)} = j\tau_0, \quad j = \text{positive integer.}$$

We will establish (1.2), (1.3), and (1.4) when a(t) satisfies (H1), (H2), and (H3) a(t) admits no representation (1.5) such that (1.6) holds.

Complementary to (H3) is

(H4) a(t) satisfies (1.5) and (1.6).

When (H4) holds, we define $\gamma = (3\delta + 2c)/(\delta + c)$ and let

$$u_1(t) = u(t) - 2\gamma^{-1} \cos \omega t$$

and

$$w_1(t) = w(t) - 2(\gamma \omega)^{-1} \sin \omega t.$$

We prove

THEOREM. Let $c \ge 0$, and let a(t) satisfy (H1) and (H2). Then

(i)
$$\lim_{t \to \infty} u(t) = 0, \quad \text{if (H3) holds,}$$
$$\lim_{t \to \infty} u_1(t) = 0, \quad \text{if (H4) holds.}$$

(ii) If
$$c+a(t) \notin L_1(0, \infty)$$
,
$$\lim_{t\to\infty} w(t) = 0, \quad \text{if (H3) holds},$$

$$\lim_{t\to\infty} w_1(t) = 0, \quad \text{if (H4) holds}.$$

(iii) If
$$c+a(t) \in L_1(0, \infty)$$
,
$$\lim_{t \to \infty} w(t) = \left(\int_0^\infty a(t) dt\right)^{-1}, \quad \text{if (H3) holds,}$$

$$\lim_{t \to \infty} w_1(t) = \left(\int_0^\infty a(t) dt\right)^{-1}, \quad \text{if (H4) holds.}$$

Generalizations of the results in [5] to nonlinear versions of (1.1) are given by Levin and Nohel in [7]. A result of Friedman (Theorem C of [3]) implies (1.2) and (1.4) for $c+a(t)=t^{-\alpha}$, $0<\alpha<1$. Halanay [4] studied a nonlinear equation

including (1.1) with k=0 when $c+a(s-\tau)-\varepsilon_0e^{-\alpha|s-\tau|}$ is a positive kernel on $\{0 \le s \le t, \ 0 \le \tau \le t\}$ for all $t \ge 0$ and some $\varepsilon_0 > 0$, $\alpha > 0$.

The Laplace transform argument of our proof resembles the proofs of the "indirect abelian" theorems in [2, pp. 265–275]. Such theorems were used by Levin and Nohel in [6] to find an asymptotic expansion as $t \to \infty$ of solutions of an equation similar to (1.1) but where the kernel is, among other things, completely monotonic on $[0, \infty)$.

Throughout the discussion S denotes the subset of the complex plane given by

$$S = \{s \mid \text{Re } s \ge 0, s \ne 0\}.$$

We define

$$A(s) = \lim_{T \to \infty} \int_0^T e^{-st} a(t) dt, \quad s \in S.$$

Then A(s) is the Laplace transform of a(t); similarly let X(s) be the Laplace transform of x(t). Taking Laplace transforms formally in (1.1), we obtain $X(s)p(s) = x_0 + (k/s)$, where p(s) = (c/s) + A(s) + s. In Lemma 5 we show that when (H3) holds, $p(s) \neq 0$ for $s \in S$. Then

$$X(s) = (x_0 + (k/s))/p(s), \quad s \in S.$$

The complex inversion formula for Laplace transforms, together with contour integration and some estimates on A(s), yields

(1.7)
$$x(t) = \lim_{\varepsilon \to 0} \frac{1}{2\pi} \left[\int_{-\infty}^{-\varepsilon} + \int_{\varepsilon}^{\infty} \right] e^{i\tau t} X(i\tau) d\tau, \qquad t > 0,$$

where for each $\varepsilon > 0$ the integrals are uniformly convergent in $t \ge T > 0$. The Riemann-Lebesgue theorem and other abelian arguments then yield our results.

When (H4) holds, we find that p(s) has exactly the two zeros $s = \pm i\omega$ in S. A formula similar to (1.7) but with principal values taken also at $\tau = \pm \omega$ is used in this case.

§2 presents a sequence of lemmas concerning a(t) and A(s). The theorem is proved in §3.

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2. The Laplace transform of a(t).

LEMMA 1. Let $\sigma > 0$, and let a(t) satisfy (H1). Then

- (i) $e^{-\sigma t}a(t)$ satisfies (H1).
- (ii) If a(t) satisfies (H2), so does $e^{-\sigma t}a(t)$.

Proof. (i) Obvious.

(ii) Since a(t) is continuous, it suffices to prove the convexity relation in (H2) with $\varepsilon = 1/2$. Set $a_i = a(t_i)$ and $b_i = \exp(-\sigma t_i)$, i = 1, 2, 3. Then a_i , $b_i \ge 0$, $a_1 - a_2 \ge a_2 - a_3 \ge 0$, $b_1 - b_2 > b_2 - b_3 > 0$. Hence $a_1b_1 - a_2b_2 = a_1(b_1 - b_2) + b_2(a_1 - a_2) \ge a_2(b_2 - b_3) + b_3(a_2 - a_3) = a_2b_2 - a_3b_3$, and (ii) is proved.

We state without proof the following easy consequence of convexity:

LEMMA 2. If a(t) satisfies (H2), then for any $\delta > 0$, the function $a(t) - a(t + \delta)$ is nonincreasing.

LEMMA 3. Let a(t) satisfy (H1). Then

- (i) A(s) is defined, finite, and continuous in S. A(s) is holomorphic in $\{\text{Re } s>0\}$.
- (ii) For $\sigma + i\tau \in S$, $\tau \neq 0$,

(2.1)
$$|\operatorname{Im} A(\sigma + i\tau)| \leq \int_0^{\pi/|\tau|} a(t) \sin |\tau| t \, dt$$
$$\leq \int_0^{\pi/|\tau|} a(t) \, dt,$$

and

$$|\operatorname{Re} A(\sigma + i\tau)| \leq \int_0^{\pi/2|\tau|} a(t) dt,$$

so that $|A(\sigma+i\tau)| \to 0$ as $|\tau| \to \infty$, uniformly in $0 \le \sigma < \infty$.

(iii) If a(t) also satisfies (H2) and $\sigma + i\tau \in S$, then

(2.3)
$$|A(\sigma + i\tau)| \ge \frac{1}{\sqrt{2}} \int_0^{\pi/2|\tau|} \cos \tau t \, e^{-\sigma t} a(t) \, dt$$

$$\ge \frac{1}{2\sqrt{2}} \int_0^{\pi/3|\tau|} e^{-\sigma t} a(t) \, dt$$

(the case $\tau=0$, $\tau^{-1}=\infty$, is included); and if $\tau>0$,

(2.4)
$$\lim_{T\to\infty}\int_0^T a(t)\sin\tau t\,dt + \lim_{T\to\infty}\int_{\pi/2\tau}^T a(t)\cos\tau t\,dt \ge 0.$$

Proof. (i) For $s = \sigma + i\tau \in S$, $\tau > 0$, T > 0, define

$$\phi(T, s) = \int_0^T a(t)e^{-\sigma t} \cos \tau t \, dt$$

$$\psi(T, s) = \int_0^T a(t)e^{-\sigma t} \sin \tau t \, dt.$$

Since $a(t) \to 0$ as $t \to \infty$, A(s) has a nonpositive abscissa of convergence (see e.g. [8, Chapter II]) so that A(s) is holomorphic in $\{\text{Re } s > 0\}$ with

(2.5)
$$[A(s)]^* = A(s^*) = \lim_{T \to \infty} [\phi(T, s) + i\psi(T, s)]$$

for Im s>0, *=complex conjugate.

For any T>0, ϕ and ψ are continuous functions of s. (H1) and Lemma 1(i) show that $e^{-\sigma t}a(t)$ is nonnegative and nonincreasing on $(0, \infty)$. For $s=\sigma+i\tau$, $\tau>0$, $T_1, T_2 \ge (n+\frac{1}{2})\pi/\tau$, n= nonnegative integer,

$$|\phi(T_1, s) - \phi(T_2, s)| \le \int_{(n+1/2)\pi/\tau}^{(n+3/2)\pi/\tau} a(t) dt$$

and similarly for ψ . Since $a(t) \to 0$ as $t \to \infty$, $\dot{\phi}(T, s)$ and $\psi(T, s)$ converge as $T \to \infty$, uniformly in any set of the form

$$S \cap \{\sigma + i\tau \mid 0 < \tau_0 \le \tau \le \tau_1 < \infty\}$$

to continuous functions $\phi(s)$ and $\psi(s)$. Comparing this with (2.5), we see that (i) is proved.

(ii) Since Re $A(\sigma+i\tau)$ and Im $A(\sigma+i\tau)$ are respectively even and odd in τ , we may assume $\tau>0$. Since for each T>0 we have

$$|\phi(T, \sigma + i\tau)| \leq \int_0^{\pi/2\tau} a(t)e^{-\sigma t} \cos \tau t \, dt$$
$$\leq \int_0^{\pi/2\tau} a(t) \, dt,$$

- (2.2) holds. (2.1) is obtained similarly.
- (iii) The case $\tau = 0$ is trivial, and by symmetry of Re $A(\sigma + i\tau)$ and Im $A(\sigma + i\tau)$ in τ , we may assume $\tau > 0$. Note first that since $A(i\tau) = \phi(i\tau) i\psi(i\tau)$,

$$(2.6) \qquad \sqrt{2|A(i\tau)|} \ge |\phi(i\tau)| + |\psi(i\tau)| \ge \phi(i\tau) + \psi(i\tau).$$

But

$$\phi(i\tau) = \int_0^{\pi/2\tau} a(t) \cos \tau t \, dt + \lim_{n \to \infty} \sum_{k=1}^n \int_{2(k-1)\pi/\tau}^{2k\pi/\tau} a\left(t + \frac{\pi}{2\tau}\right) \cos \left[\tau\left(t + \frac{\pi}{2\tau}\right)\right] \, dt$$

and

$$\psi(i\tau) = \lim_{n \to \infty} \sum_{k=1}^{n} \int_{2(k-1)\pi/\pi}^{2k\pi/\tau} a(t) \sin \tau t \, dt,$$

so (2.6) becomes

$$|\sqrt{2}A(i\tau)| - \int_0^{\pi/2\tau} a(t) \cos \tau t \, dt$$

$$\geq \lim_{n \to \infty} \sum_{k=1}^n \int_{2(k-1)\pi/\tau}^{2k\pi/\tau} \left\{ a(t) \sin \tau t + a\left(t + \frac{\pi}{2\tau}\right) \cos\left[\tau\left(t + \frac{\pi}{2\tau}\right)\right] \right\} dt.$$

We note that the right-hand side of (2.7) equals the left-hand side of (2.4) and that both (2.3) for $\sigma = 0$ and (2.4) will follow if we show that this right-hand side is nonnegative. But for any integer $k \ge 1$,

$$\int_{2(k-1)\pi/\tau}^{2k\pi/\tau} \left\{ a(t) \sin \tau t + a\left(t + \frac{\pi}{2\tau}\right) \cos \left[\tau\left(t + \frac{\pi}{2\tau}\right)\right] \right\} dt$$

$$= \int_{0}^{\pi/\tau} \sin \tau t \left[a(t+x_0) - a(t+x_1) - a(t+x_2) + a(t+x_3)\right] dt,$$

where $x_j = 2(k-1)\pi/\tau + j\pi/2\tau$, j=0, 1, 2, 3, and Lemma 2 with $\delta = \pi/2\tau$ shows that the integrand is nonnegative. For (2.3) with $\sigma > 0$, we apply (2.3) with $\sigma = 0$ to the function $b(t) = e^{-\sigma t}a(t)$, which satisfies (H1) and (H2) by Lemma 1, and which has Laplace transform $B(s) = A(s+\sigma)$. This completes the proof of Lemma 3.

COROLLARY 3.1. Let a(t) satisfy (H1). Then $|sA(s)| \to 0$ as $s \to 0$, $s \in S$.

Proof. We let $s = \sigma + i\tau$. Applying (2.1) and (2.2) to the function $b(t) = e^{-\sigma t}a(t)$, we have

$$|sA(s)| = |s| |B(i\tau)| \leq \sqrt{2}|s| \int_0^{\pi/|\tau|} e^{-\sigma t} a(t) dt,$$

when $\tau \neq 0$ and, trivially, also when $\tau = 0$, $\sigma > 0$. Thus if $\sigma \geq |\tau|$, $|sA(s)| \leq 2\sigma \int_0^\infty e^{-\sigma t} \cdot a(t) dt = 2\sigma A(\sigma)$, while if $|\tau| > \sigma$, we have $|sA(s)| \leq 2|\tau| \int_0^{\pi/|\tau|} a(t) dt$; one of these estimates is valid for each $s \in S$. But since $a(t) \to 0$ as $t \to \infty$,

(2.8)
$$\int_0^x a(t) dt = o(x) \qquad (x \to \infty).$$

It follows that $|\tau| \int_0^{\pi/|\tau|} a(t) dt \to 0$ as $\tau \to 0$; and (2.8) together with an elementary abelian theorem for Laplace transforms [8, p. 182] gives $\sigma A(\sigma) \to 0$ as $\sigma \to 0+$. In view of our estimates for |sA(s)|, the corollary is proved.

COROLLARY 3.2. Let a(t) satisfy (H1) and (H2), and suppose that $a(t) \notin L_1(0, \infty)$. Then $[s+A(s)]^{-1} \to 0$ as $s \to 0$, $s \in S$.

Proof. By Lemma 3(iii),

$$|A(s)| = |A(\sigma + i\tau)| \ge \frac{1}{2\sqrt{2}} \int_0^{\pi/3|\tau|} e^{-\sigma t} a(t) dt$$
$$\ge m \int_0^{\pi/3|s|} a(t) dt,$$

where $m^{-1} = 2\sqrt{2}e^{\pi/3}$. Then for sufficiently small |s|,

$$|[s+A(s)]^{-1}| \leq 2\Big(m\int_0^{\pi/3|s|} a(t) dt\Big)^{-1} = o(1) \qquad (|s| \to 0).$$

LEMMA 4. Suppose a(t) satisfies (H1) and (H2). Then exactly one of the following two cases holds:

I. Either (i) $a(0+) = \infty$ or (ii) $a(0) \equiv a(0+) < \infty$ and $\forall \tau > 0$ there exists an integer $k = k(\tau) > 0$ such that

$$a\left(\frac{2(k-1)\pi}{\tau}\right) - 2a\left(\frac{2(k-1)\pi + \pi}{\tau}\right) + a\left(\frac{2k\pi}{\tau}\right) > 0.$$

II. (i) There exists a positive number τ_0 and a sequence $\{\delta_k\}_{k=1}^{\infty}$ such that (1.5) holds.

(ii) The numbers τ_0 , t_0 , and δ and the sequence $\{\delta_k\}$ are determined uniquely by (1.5). All positive τ such that

$$(2.10) a\left(\frac{2(k-1)\pi}{\tau}\right) - 2a\left(\frac{2(k-1)\pi + \pi}{\tau}\right) + a\left(\frac{2k\pi}{\tau}\right) = 0, k = 1, 2, \dots$$

are integral multiples of τ_0 .

(iii) The Laplace transform of a(t) is given by

(2.11)
$$A(s) = \frac{\delta}{s} + \frac{1}{s^2} \sum_{k=1}^{\infty} \frac{\delta_k}{kt_0} (\exp[-skt_0] - 1), \quad s \in S.$$

Proof. First we note that a representation (1.5) implies (2.10) with $\tau = \tau_0$, so cases I and II exclude each other.

Now suppose we are not in case I, so that (2.10) holds for some $\tau = \tau_1$. Let J denote the set of positive integers j such that (2.10) holds when $\tau = \tau_1/j$. Then $1 \in J$. Also, J is a finite set; for (H1), (H2), and (2.10) with k = 1, $\tau = \tau_1/j$, $j \in J$, show that a(t) is linear with negative slope on $[0, 2j\pi/\tau_1]$ whenever $j \in J$. We let j_0 be the largest $j \in J$, and set $\tau_0 = \tau_1/j_0$. Then for any integer j > 1, (2.10) does *not* hold with $\tau = \tau_0/j$.

By convexity, (2.10) with $\tau = \tau_0$ shows that a(t) is linear on each interval $2(k-1)\pi/\tau_0 \le t \le 2k\pi/\tau_0$; we let $-\lambda_k$ be its slope there. By (H1) and (H2),

$$\lambda_1 \ge \lambda_2 \ge \lambda_3 \ge \cdots \ge 0$$
, and $\lim_{k \to \infty} \lambda_k = 0$.

We define $t_0 = 2\pi/\tau_0$ and $\delta_k = kt_0(\lambda_k - \lambda_{k+1}) \ge 0$, $k = 1, 2, 3, \ldots$. Then on the interval $(k-1)t_0 \le t \le kt_0$, the function defined by the right-hand side of (1.5a) has the value

$$\sum_{j=k}^{\infty} \delta_j \left(1 - \frac{t}{jt_0} \right) = \sum_{j=k}^{\infty} (\lambda_j - \lambda_{j+1}) (jt_0 - t)$$

$$= \lambda_k (kt_0 - t) + \sum_{j=k+1}^{\infty} \lambda_j t_0$$

$$= \int_0^t da(\tau) = a(t).$$

This proves (1.5a) and (1.5b).

For (1.5c), we note that if j > 1 divides all k in Ω , then

$$a(t) = \sum_{k=1}^{\infty} \delta_{jk} \left(1 - \frac{\min\{t, jkt_0\}}{jkt_0} \right),$$

and as in the proof of (1.5a), a(t) is linear on $2j(k-1)\pi/\tau_0 \le t \le 2jk\pi/\tau_0$, $k=1, 2, \ldots$. But then (2.10) holds with $\tau = \tau_0/j$, and we chose τ_0 so as to make this impossible. This completes the proof that II(i) holds when I does not hold.

To prove (ii), suppose

(1.5a')
$$a(t) = \sum_{k=1}^{\infty} \delta_k' \left(1 - \frac{\min\{t, kt_0'\}}{kt_0'} \right), \quad t_0' = 2\pi/\tau_0'$$

with corresponding (1.5b'), (1.5c'). Let $k_1 < k_2 < k_3 < \cdots$ be all the elements of Ω , and $k'_1 < k'_2 < \cdots$ all the elements of Ω' . By (1.5a, a'), for each i

(2.12) $k_i t_0 = k'_i t'_0 = \max\{x \mid \text{slope of } a(t) \text{ has exactly } i \text{ different values on } (0, x)\}.$

In particular, $t_0/t_0' = k_1'/k_1$ = rational number, so $t_0/t_0' = p/q$, where p and q are relatively prime positive integers. Then for each i, by (2.12), $k_i p/q = k_i' = \text{integer}$. By (1.5c), q = 1, so by (1.5c') also p = 1 and $t_0 = t_0'$. By (1.5a, a'), $\tau_0 = \tau_0'$. By (1.5a'), the slope $-\lambda_k$ of a(t) in $[(k-1)t_0, kt_0]$ is $-\sum_{j=k}^{\infty} (\delta_j'/jt_0)$, so by the definition of δ_k in the proof of (i), $\delta_k' = (\lambda_k - \lambda_{k+1})kt_0 = \delta_k$. (1.5b, b') give $\delta = \delta'$, and uniqueness is proved.

Any τ satisfying (2.10) leads, as in (i), to a representation (1.5') with $\tau'_0 = \tau/j$. By uniqueness $j\tau_0 = j\tau'_0 = \tau$, and (ii) is proved.

(iii) This follows from (1.5) by direct computation. This completes the proof of Lemma 4.

LEMMA 5. Let a(t) satisfy (H1) and (H2), and let $c \ge 0$. Define

$$p(s) = (c/s) + A(s) + s, \qquad s \in S.$$

Then

- (i) p(s) has no zeros in S if (H3) holds. If (H4) holds p(s) has exactly the two zeros $s = \pm i\omega$ in S.
 - (ii) When (H4) holds,

$$(2.13) |p(s)-\gamma(s-i\omega)| = o(|s-i\omega|) (s \to i\omega, s \in S),$$

where $\gamma = (3\delta + 2c)/(\delta + c)$.

Proof. (i) First, (2.11) shows that $p(\pm i\omega) = 0$ when (H4) holds. For $s = i\tau \neq 0$,

$$\operatorname{Re} A(i\tau) = \lim_{n \to \infty} \sum_{k=1}^{n} \int_{0}^{2\pi/|\tau|} a\left(\frac{2(k-1)\pi}{|\tau|} + t\right) \cos \tau t \, dt.$$

But for each k, the integral in the sum is equal to

$$\int_0^{\pi/2|\tau|} \left[a(x_0+t) - a(x_1-t) - a(x_1+t) + a(x_2-t) \right] \cos \tau t \, dt,$$

where $x_j = [2(k-1)+j]\pi/|\tau|$, j=0, 1, 2. Lemma 2 with $\delta = (\pi/|\tau|)-2t$ shows that the integrand is nonnegative. Furthermore, Lemma 4 shows that if a(t) is in case I of Lemma 4 or in case II with $j\tau_0 \neq |\tau|$, $j=1, 2, 3, \ldots$, then there exists k such that the integrand is positive at t=0, and by continuity on an interval $0 \leq t < \varepsilon$; for this k the integral is positive. We conclude that Re $A(i\tau) \geq 0$, and if Re $A(i\tau) = 0$, then a(t) is in case II of Lemma 4 with $\tau = j\tau_0$, j = integer.

By Lemma 1, if $\sigma > 0$, the function $e^{-\sigma t}a(t)$ satisfies (H1) and (H2). Thus for $\sigma > 0$ (by the preceding paragraph for $\tau \neq 0$ and trivially for $\tau = 0$), Re $A(\sigma + i\tau) = \int_0^\infty \left[e^{-\sigma t}a(t) \right] \cos \tau t \ dt \ge 0$.

To apply these remarks, we suppose that p(s) = 0, $s = \sigma + i\tau \in S$. Then

$$0 = \operatorname{Re} p(s) = c\sigma/(\sigma^2 + \tau^2) + \operatorname{Re} A(\sigma + i\tau) + \sigma.$$

Therefore $\sigma = 0$ and a(t) is in case II with $\tau = j\tau_0$, j = integer. But then, using (2.11),

$$0 = \text{Im } p(s) = -(c+\delta)/j\tau_0 + j\tau_0;$$

i.e., $c + \delta = (j\tau_0)^2 = \tau^2$, so that (H4) holds and $s = \pm i\omega$. This proves (i). (ii) By (2.11),

$$p(s)-s+\frac{\omega^2}{s}+\frac{\delta}{s^2}(s-i\omega)=\left[\frac{1}{s^2}\sum_{k=1}^{\infty}\frac{\delta_k}{kt_0}\left(\exp\left[-skt_0\right]-1\right)\right]-\frac{\delta}{s^2}(s-i\omega).$$

On the left-hand side we expand s, ω^2/s , and δ/s^2 in Taylor series about $i\omega$; on the right we rearrange terms. Then

(2.14)
$$p(s) - \gamma(s - i\omega) + O(|s - i\omega|^2) = s^{-2} \sum_{k=1}^{\infty} \frac{\delta_k}{kt_0} (\exp[-skt_0] - 1 + skt_0 - i\omega kt_0)$$
 $(s \to i\omega, s \in S).$

For Re $s \ge 0$,

$$|(\exp[-skt_0] - 1 + skt_0 - i\omega kt_0)/kt_0| \le |\exp[-skt_0] - 1|/kt_0 + |s - i\omega|$$
$$\le 2/kt_0 + |s - i\omega|.$$

On the other hand, the power series expansion of $\exp(-skt_0)$ about $s=i\omega$ yields for $k|s-i\omega| \le 1$

$$\begin{aligned} |(\exp[-skt_0] - 1 + skt_0 - i\omega kt_0)/kt_0| &= \left| \sum_{j=2}^{\infty} \left[(s - i\omega)(-kt_0) \right]^j / j! \right| / kt_0 \\ &= k|s - i\omega|^2 \left| \sum_{j=0}^{\infty} t_0^{j+2} / (j+2)! \right| / t_0 \\ &\leq (\exp[t_0]k/t_0) |s - i\omega|^2. \end{aligned}$$

Using these two estimates, we have for Re $s \ge 0$

(2.15)
$$\left| \sum_{k=1}^{\infty} \frac{\delta_{k}}{kt_{0}} \left(\exp\left[-skt_{0} \right] - 1 + skt_{0} - i\omega kt_{0} \right) \right|$$

$$\leq \frac{\exp\left[t_{0} \right]}{t_{0}} |s - i\omega|^{2} \sum_{k=1}^{n(s)} k\delta_{k} + |s - i\omega| \sum_{k=n(s)+1}^{\infty} \delta_{k} + \frac{2}{t_{0}} \sum_{k=n(s)+1}^{\infty} \delta_{k}/k,$$

where n(s) is the greatest integer such that $n(s)|s-i\omega| \le 1$.

By (1.5a), $a(t) \ge \sum_{k=1}^{m} \delta_k [1 - (\min\{t, kt_0\})/kt_0]$ for $m = 1, 2, 3, \dots$ Since

$$\int_0^{kt_0} \delta_k(1-t/kt_0) dt = kt_0\delta_k/2,$$

we have $2 \int_0^{mt_0} a(t) dt \ge t_0 \sum_{k=1}^m k \delta_k$. Hence

$$(2.16) |s-i\omega| \sum_{k=1}^{n(s)} k \delta_k \le \frac{2}{t_0 n(s)} \int_0^{t_0 n(s)} a(t) dt \to 0 as |s-i\omega| \to 0,$$

since $a(t) \to 0$. Also, since $\sum \delta_k < \infty$,

$$\sum_{k=n(s)+1}^{\infty} \delta_k + \frac{2}{t_0|s-i\omega|} \sum_{k=n(s)+1}^{\infty} \delta_k/k \le (1+2/t_0) \sum_{k=n(s)+1}^{\infty} \delta_k \to 0 \quad \text{as } |s-i\omega| \to 0.$$

Combining this with (2.14), (2.15), and (2.16), we have (2.13), and Lemma 5 is proved.

3. **Proof of theorem.** Integrating (1.1), we obtain

$$x(t) = x_0 + kt - \int_0^t f(t-\tau)x(\tau) d\tau,$$

where $0 \le f(t) = \int_0^t [a(\tau) + c] d\tau \le \int_0^1 a(\tau) d\tau + c + [a(1) + c]t$. By a standard result on Volterra equations [1, §7.6], x(t) satisfies an inequality

$$|x(t)| \leq B_1 e^{bt}, \quad b, B_1 > 0.$$

Substituting in (1.1), we have

$$|x'(t)| \leq k + B_1 e^{bt} \int_0^1 [a(\tau) + c] d\tau + [a(1) + c] \int_0^t B_1 e^{b\tau} d\tau \leq B_2 e^{bt}.$$

Taking Laplace transforms in (1.1), we obtain $X(s)p(s) = x_0 + (k/s)$, Re s > b, with p(s) = c/s + A(s) + s, as in Lemma 5. By Lemma 5(i),

(3.1)
$$X(s) = (x_0 + k/s)/p(s), \quad \text{Re } s > b,$$

and (3.1) defines X(s) as a function holomorphic in $\{\text{Re } s > 0\}$ and continuous in S (for H3) or in $S - \{\pm i\omega\}$ (for H4). Also note that by (2.5) we have $[X(s)]^* = X(s^*)$; and by Lemma 3(ii), $X(\sigma + i\tau) \to 0$ as $|\tau| \to \infty$, uniformly in $0 \le \sigma < \infty$.

(i) We set $x_0 = 1$, k = 0, X = U in (3.1). Then

$$U(s) = \frac{1}{s} - \frac{(c/s) + A(s)}{c + sA(s) + s^2}, \qquad s = \sigma + i\tau \in S.$$

For any $\sigma \ge 0$ and sufficiently large R > 0, the second term is in $L_1\{(-\infty, -R) \cup (R, \infty)\}$ as a function of τ , by Lemma 3(ii); and integration by parts shows that for any T > 0

$$\left[\int_{-\infty}^{-R} + \int_{R}^{\infty}\right] e^{i\tau t} (\sigma + i\tau)^{-1} d\tau$$

converges uniformly for $t \ge T$. Then the exponential bound on u(t) and $u(t) \in C'$ justify the inversion formula

(3.2)
$$2\pi u(t) = e^{\sigma t} \int_{-\infty}^{\infty} e^{i\tau t} U(\sigma + i\tau) d\tau, \qquad \sigma > b, t > 0.$$

If $c+a(t) \in L_1(0, \infty)$, A(s) has limit $A(0) = \int_0^\infty a(t) dt$ at s=0, so U(s) is continuous with U(0) = 1/A(0). If $c+a(t) \notin L_1(0, \infty)$, Corollaries 3.1 (if c>0) and 3.2 (if c=0) show that $U(s) \to 0$ as $s \to 0$, $s \in S$, and again U(s) is continuous at s=0.

Thus if (H3) holds, Cauchy's theorem and the fact that $U(\sigma+i\tau)\to 0$ as $|\tau|\to\infty$ uniformly in σ yield

(3.3)
$$2\pi u(t) = \int_{-\infty}^{\infty} e^{i\tau t} U(i\tau) d\tau.$$

The Riemann-Lebesgue theorem for finite intervals and the uniform convergence of the integral in (3.3) yield $u(t) \to 0$ as $t \to \infty$.

On the other hand, if (H4) holds, we set

(3.4)
$$U_1(s) = U(s) - 2s/\gamma(s^2 + \omega^2).$$

Since $2s/\gamma(s^2+\omega^2)$ is the Laplace transform of $2\gamma^{-1}\cos\omega t$, (3.2) holds with u_1 and U_1 in place of u and U. Using Cauchy's theorem as before, we have for $0 < \rho < \omega$,

$$2\pi u_1(t) = \left[\int_{-\infty}^{-\omega-\rho} + \int_{-\omega+\rho}^{\omega-\rho} + \int_{\omega+\rho}^{\infty}\right] e^{i\tau t} U_1(i\tau) d\tau$$
$$+ \frac{1}{i} \left[\int_{C_0^-} + \int_{C_0^+} e^{st} U_1(s) ds, \quad t > 0,\right]$$

where C_{ρ}^{\pm} is the semicircle $\{\pm i\omega + \rho e^{i\theta}, -\pi/2 \le \theta \le \pi/2\}$. Since $U_1(s^*) = [U_1(s)]^*$, this may also be written

$$(3.5) \quad \pi u_1(t) = \left[\int_0^{\omega - \rho} + \int_{\omega + \rho}^{\infty} \right] \operatorname{Re} \left\{ e^{i\tau t} U_1(i\tau) \right\} d\tau + \frac{1}{i} \int_{C_{\sigma}^+} \operatorname{Re} \left\{ e^{st} U_1(s) \right\} ds, \quad t > 0.$$

Note that

$$U_1(s) = \frac{1}{p(s)} - \frac{1}{\gamma(s-i\omega)} + O(1) \qquad (s \to i\omega, s \in S).$$

Writing $p(s) = \gamma(s - i\omega) + [p(s) - \gamma(s - i\omega)]$, we find that

$$(3.6) \quad U_1(s) = \frac{p(s) - \gamma(s - i\omega)}{\gamma(s - i\omega)^2} \left[\frac{-1}{\gamma + (p(s) - \gamma(s - i\omega))/(s - i\omega)} \right] + O(1) \quad (s \to i\omega, s \in S).$$

Thus, for $s \in C_{\rho}^+$, (2.13) yields $e^{st}U_1(s) = o(\rho^{-1})$, $(\rho \to 0)$. Since $|C_{\rho}^+| = \pi \rho$, we may let $\rho \to 0$ in (3.5) and obtain for $0 < \eta < \omega$

$$\pi u_1(t) = \left[\int_0^{\omega - \eta} + \int_{\omega + \eta}^{\infty} \right] \operatorname{Re} \left\{ e^{i\tau t} U_1(i\tau) \right\} d\tau + \lim_{\varepsilon \to 0} \left[\int_{\omega - \eta}^{\omega - \varepsilon} + \int_{\omega + \varepsilon}^{\omega + \eta} \right] \operatorname{Re} \left\{ e^{i\tau t} U_1(i\tau) \right\} d\tau.$$

Treating the first term as for (H3) (3.4 and integration by parts show that $\int_{\omega+\eta}^{\infty}$ converges uniformly), we have

(3.7)
$$\pi u_1(t) = \lim_{\varepsilon \to 0} \left[\int_{\omega - \eta}^{\omega - \varepsilon} + \int_{\omega + \varepsilon}^{\omega + \eta} \right] \{ [\text{Re } U_1(i\tau)] \cos \tau t - [\text{Im } U_1(i\tau)] \sin \tau t \} d\tau + o(1),$$

$$(t \to \infty), \ 0 < \eta < \omega.$$

For real λ define

$$S(\lambda) = \sum_{k=1}^{\infty} \frac{\delta_k}{kt_0} (\sin kt_0 \lambda - kt_0 \lambda), \qquad C(\lambda) = \sum_{k=1}^{\infty} \frac{\delta_k}{kt_0} (\cos kt_0 \lambda - 1).$$

Note that

(3.8)
$$C(\lambda) = C(-\lambda) \le 0, \quad \lambda S(\lambda) = -\lambda S(-\lambda) \le 0,$$

and $C(\lambda + j\omega) = C(\lambda)$, for any integer j. By (2.14),

$$p(i\tau) - i\gamma(\tau - \omega) = \left[-C(\tau) + iS(\tau - \omega) \right] / \tau^2 + O(|\tau - \omega|^2) \qquad (\tau \to \omega).$$

(2.15) and the argument following it show that

(3.9)
$$|C(\lambda)| + |S(\lambda)| = o(\lambda), \qquad \lambda \to 0.$$

Using these facts in (3.6) one computes

(3.10)
$$\operatorname{Re} U_1(i\tau) = \frac{-C(\tau)}{(\tau - \omega)^2 \tau^2 [\gamma^2 + o(1)]} + O(1) \qquad (\tau \to \omega),$$

and

(3.11)
$$\operatorname{Im} U_1(i\tau) = R(\tau - \omega) + o\left(\frac{C(\tau)}{(\tau - \omega)^2}\right) + O(1) \qquad (\tau \to \omega),$$

where

$$R(\lambda) = \frac{S(\lambda)[\gamma + S(\lambda)/\omega^2 \lambda]}{\gamma \lambda^2 \omega^2 [(\gamma + S(\lambda)/\lambda \omega^2)^2 + C^2(\lambda)/\lambda^2 \omega^4]}$$

Now let $t=t^*=2\pi/\omega$ in (3.7). Since $\sin \tau t^*=\sin (\tau-\omega)t^*=O(\tau-\omega)$, $(\tau\to\omega)$, we see from (3.9) and (3.11) that with $t=t^*$ the second term in the integrand in (3.7) is bounded on $(\omega-\eta,\omega+\eta)$. It follows from (3.7) and (3.10) that

(3.12)
$$\lim_{\varepsilon \to 0} \left[\int_{\omega - \tau}^{\omega - \varepsilon} + \int_{\omega + \varepsilon}^{\omega + \eta} \left[\frac{C(\tau) \cos \tau t^* d\tau}{(\tau - \omega)^2 \tau^2 [\gamma^2 + o(1)]} \right] \right]$$

exists and is finite. But by the choice of t^* , the integrand in (3.12) is $\leq C(\tau)/2(\tau-\omega)^2\gamma^2\omega^2$ for $|\tau-\omega|$ sufficiently small. Since $C(\tau)\leq 0$ we conclude

(3.13)
$$C(\tau)/(\tau-\omega)^2 \in L_1(\omega-\eta,\,\omega+\eta).$$

In view of (3.10), (3.11), and (3.13), an application of the Riemann-Lebesgue theorem to (3.7) yields

$$(3.14) \quad -\pi u_1(t) = \lim_{\varepsilon \to 0} \left[\int_{\omega-\tau}^{\omega-\varepsilon} + \int_{\omega+\varepsilon}^{\omega+\eta} R(\tau-\omega) \sin \tau t \, d\tau + o(1), \quad t \to \infty. \right]$$

Note that $R(-\lambda) = -R(\lambda)$. The change of variables $\lambda = \tau - \omega$ in (3.14) shows that to complete the proof of (i) we need only show that

$$(3.15) \quad r(t) \equiv \lim_{\varepsilon \to 0} \int_{s}^{\eta} \left[\sin(\omega + \lambda)t - \sin(\omega - \lambda)t \right] R(\lambda) \, d\lambda = o(1) \qquad (t \to \infty).$$

A trigonometric identity and $|(\sin \lambda t)/\lambda| \le t$ permit us to write

(3.16)
$$r(t) = 2\cos(\omega t) \int_{0}^{\pi} R(\lambda) \sin \lambda t \, d\lambda.$$

To prove (3.15), note first that

$$|S'(\lambda)| = \left|\sum_{k=1}^{\infty} \delta_k(\cos kt_0\lambda - 1)\right| \le 2\delta.$$

Similarly, $|C'(\lambda)| \le \delta$. Straightforward computations and estimates then show that $|R'(\lambda)| \le M/\lambda^2$ for some constant $M < \infty$, $0 < \lambda < \eta$. Also note that $|\lambda R(\lambda)| \le K_1 |S(\lambda)/\lambda| \le K_2$, K_1 , $K_2 < \infty$.

Now let $\varepsilon > 0$ be given; pick $\mu > 0$ so that $(M + K_2) < \mu \varepsilon / 3$, and pick T > 0 so that $(M + K_2) / \eta T \le \varepsilon / 3$ and $|K_1 S(\lambda)/\lambda| \le \varepsilon / 3\mu$ for $0 < \lambda < \mu / T$. Then for $t \ge T$, integration by parts yields

$$\left| \int_{\mu/t}^{\eta} R(\lambda) \sin \lambda t \ d\lambda \right| \leq (M + K_2) \left(\frac{1}{\eta T} + \frac{1}{\mu} \right) \leq 2\varepsilon/3,$$

while

$$\left| \int_0^{\mu/t} R(\lambda) \sin \lambda t \, d\lambda \right| \le \left| \int_0^{\mu/t} \left[\lambda R(\lambda) \sin \lambda t \right] / \lambda \, d\lambda \right| \le \varepsilon/3.$$

These estimates, together with (3.16), prove (3.15) and complete the proof of (i). (ii) We set $x_0 = 0$, k = 1, X = W in (3.1). Note that $W(\sigma + i\tau) = O(\tau^{-2})$, $|\tau| \to \infty$, so $W(\sigma + i\tau) \in L_1\{(-\infty, R) \cup (R, \infty)\}$ for R sufficiently large and $0 \le \sigma < \infty$.

If (H3) holds, we use Cauchy's theorem as in (i) to obtain, for each $\varepsilon > 0$ sufficiently small,

$$(3.17) 2\pi w(t) = \left[\int_{-\infty}^{-\varepsilon} + \int_{\varepsilon}^{\infty} e^{i\tau t} W(i\tau) d\tau + \frac{1}{i} \int_{\Omega_{\varepsilon}} e^{st} W(s) ds,\right]$$

where D_{ε} is the semicircle $\{\varepsilon e^{i\theta}, -\pi/2 \le \theta \le \pi/2\}$. Corollaries 3.1 and 3.2 show that $W(\varepsilon e^{i\theta}) = o(1/\varepsilon)$, $(\varepsilon \to 0)$, uniformly in $-\pi/2 \le \theta \le \pi/2$. Since $|D_{\varepsilon}| = \varepsilon \pi$, we may let $\varepsilon \to 0$ in (3.17) and obtain

$$2\pi w(t) = \lim_{\varepsilon \to 0} \left[\int_{-\infty}^{-\varepsilon} + \int_{0}^{\infty} e^{i\pi t} W(i\tau) d\tau, \quad t > 0. \right]$$

Now the symmetry of $W(i\tau)$ in τ and the Riemann-Lebesgue theorem imply that for $\Delta > 0$,

(3.18)
$$\pi w(t) = \lim_{\varepsilon \to 0} \int_{\varepsilon}^{\Delta} \operatorname{Re} \left\{ e^{i\tau t} W(i\tau) \right\} d\tau + o(1) \qquad (t \to \infty).$$

Next we derive a formula similar to (3.18) for $w_1(t)$ in the (H4) case. Define

$$W_1(s) = W(s) - 2/\gamma(s^2 + \omega^2).$$

Note that $2/\gamma(s^2+\omega^2)$ is the Laplace transform of $2(\gamma\omega)^{-1}\sin\omega t$. Note also that $W_1(s)=U_1(s)/s=(U_1(s)/\omega)+O(1),\ (s\to\omega,\,s\in S)$. Continuing as with $U_1(s)$ in (i) and using $W_1(i\tau)\in L_1(\omega+\eta,\,\infty)$, we obtain

$$\pi w_1(t) = \lim_{\varepsilon \to 0} \int_{\varepsilon}^{\omega - \eta} \operatorname{Re} \left\{ e^{t\tau t} W_1(i\tau) \right\} d\tau$$

$$+ \omega^{-1} \lim_{\rho \to 0} \left[\int_{\omega - \eta}^{\omega - \rho} + \int_{\omega + \rho}^{\omega + \eta} \right] \operatorname{Re} \left\{ e^{t\tau t} U_1(i\tau) \right\} d\tau + o(1),$$

$$t \to \infty, \ 0 < \eta < \omega.$$

Since $W_1(s) - W(s)$ is bounded on $(0, \omega - \eta)$, we may replace W_1 by W in the first term above. By (3.7) and conclusion (i) of the theorem, the second term is o(1), $(t \to \infty)$. Thus for $0 < \Delta < \omega$,

(3.19)
$$\pi w_1(t) = \lim_{\varepsilon \to 0} \int_{\varepsilon}^{\Delta} \operatorname{Re} \left\{ e^{i\tau t} W_1(i\tau) \right\} d\tau + o(1) \qquad (t \to \infty).$$

Now set $\Delta = 1$ when (H3) holds, $\Delta = \omega/2$ when (H4) holds. By (3.18) and (3.19) we can complete the proof by showing that

(3.20)
$$\lim_{t\to\infty} \left(\lim_{\varepsilon\to\infty} \int_{\varepsilon}^{\Delta} \operatorname{Re} \left\{ e^{i\tau t} W(i\tau) \right\} d\tau \right) = 0.$$

If c>0, $W(i\tau)$ is continuous on $[0, \Delta]$, by Corollary 3.1, so that (3.20) holds. It remains to prove (3.20) with c=0.

For $0 < \tau \leq \Delta$, define

(a)
$$\phi_1(i\tau) = \int_0^{\pi/2\tau} a(t) \cos \tau t \, dt$$

(3.21) (b)
$$\phi_2(i\tau) = \lim_{T \to \infty} \int_{\pi/2\tau}^T a(t) \cos \tau t \, dt$$

(c)
$$\phi(i\tau) = \phi_1(i\tau) + \phi_2(i\tau)$$

(d)
$$\psi(i\tau) = \lim_{T\to\infty} \int_0^T a(t) \sin \tau t \, dt$$
.

As in the proof of Lemma 3,

(3.22) Re
$$A(i\tau) = \phi(i\tau)$$
 and Im $A(i\tau) = -\psi(i\tau)$.

Lemma 3 gives some useful facts about these functions. In addition, since $0 \le a(t) \downarrow$, $\phi_2(i\tau) \le 0$, so that (2.4) implies

$$(3.23) \psi(i\tau) \ge |\phi_2(i\tau)|, 0 < \tau \le \Delta.$$

Also, (2.1), (3.22), and $0 \le a(t) \downarrow \text{ yield}$

$$0 \le \psi(i\tau) \le \int_0^{\pi/\tau} a(t) \sin \tau t \, dt$$
$$\le 4 \int_0^{\pi/4\tau} a(t) \cos \tau t \, dt$$
$$\le 4 \int_0^{\pi/2\tau} a(t) \cos \tau t \, dt,$$

i.e.,

$$(3.24) 0 \leq \psi(i\tau) \leq 4\phi_1(i\tau), 0 < \tau \leq \Delta.$$

The choice of Δ insures that $W(i\tau)$ is continuous on $(0, \Delta]$. Using (3.22) and our assumption c=0, we compute

$$W(i\tau) = \frac{(\psi(i\tau) - \tau)}{\tau |A(i\tau) + i\tau|^2} - i \frac{\phi(i\tau)}{\tau |A(i\tau) + i\tau|^2}.$$

Defining $y(t) = \lim_{\varepsilon \to 0} \int_{\varepsilon}^{\Delta} \operatorname{Re} \{e^{i\tau t} W(i\tau)\} d\tau$, we have

$$y(t) = \lim_{\varepsilon \to 0} \left[\int_{\varepsilon}^{\Delta} \cos \tau t \, \frac{\psi(i\tau) \, d\tau}{\tau |A(i\tau) + i\tau|^2} - \int_{\varepsilon}^{\Delta} \frac{\cos \tau t \, d\tau}{|A(i\tau) + i\tau|^2} + \int_{\varepsilon}^{\Delta} \frac{\sin \tau t}{\tau} \, \frac{\phi(i\tau) d\tau}{|A(i\tau) + i\tau|^2} \right].$$

By Corollary 3.2, the middle integrand is continuous on $[0, \Delta]$. Since $|\sin \tau t/\tau| \le t$ for t > 0 and

$$\frac{|\phi(i\tau)|}{|A(i\tau)+i\tau|} = \frac{|\phi(i\tau)|}{|\phi(i\tau)+i(\tau-\psi(i\tau))|} \le 1,$$

Corollary (3.2) also shows that the third integrand is continuous in $[0, \Delta]$ for each t. Therefore

(3.25)
$$y(t) = \lim_{\varepsilon \to 0} \int_{\varepsilon}^{\Delta} \cos \tau t \, \frac{\psi(i\tau) \, d\tau}{\tau |A(i\tau) + i\tau|^2} - \int_{0}^{\Delta} \frac{\cos \tau t \, d\tau}{|A(i\tau) + i\tau|^2} + \int_{0}^{\Delta} \frac{\sin \tau t}{\tau} \, \frac{\phi(i\tau) \, d\tau}{|A(i\tau) + i\tau|^2}, \qquad t > 0.$$

But for $0 < \tau < \pi/3t$ we have $\cos \tau t > 1/2$ and $\psi(i\tau) \ge 0$; hence the existence of the limit in (3.25) shows that

(3.26)
$$\psi(i\tau)/\tau |A(i\tau)+i\tau|^2 \in L_1(0, \Delta).$$

Applying the Riemann-Lebesgue theorem to the first two integrals in (3.25), we obtain

$$(3.27) y(t) = \int_0^\Delta \frac{\sin \tau t}{\tau} \frac{\phi(i\tau) d\tau}{|A(i\tau) + i\tau|^2} + o(1) (t \to \infty).$$

Another consequence of (3.26), together with (3.23), is

(3.28)
$$\phi_2(i\tau)/\tau |A(i\tau)+i\tau|^2 \in L_1(0, \Delta).$$

We rewrite (3.27) as

(3.29)
$$y(t) = \int_0^\Delta \frac{\sin \tau t}{\tau} \frac{d\tau}{\phi_1(i\tau)} + \int_0^\Delta \frac{\sin \tau t}{\tau} \frac{\phi_1(i\tau)\phi(i\tau) - \phi^2(i\tau) - (\psi(i\tau) - \tau)^2}{\phi_1(i\tau)|A(i\tau) + i\tau|^2} d\tau + o(1)$$
$$= y_1(t) + y_2(t) + o(1) \qquad (t \to \infty).$$

Then

$$(3.30) \quad y_2(t) = -\int_0^\Delta \sin \tau t \left[\frac{\phi_1(i\tau) + \phi_2(i\tau)}{\phi_1(i\tau)} \phi_2(i\tau) + \frac{\psi(i\tau) - \tau}{\phi_1(i\tau)} (\psi(i\tau) - \tau) \right] \frac{d\tau}{\tau |A(i\tau) + i\tau|^2}$$

Now, (3.23) and (3.24) show that $|(\phi_1(i\tau)+\phi_2(i\tau))/\phi_1(i\tau)|$ and $|\psi(i\tau)/\phi_1(i\tau)|$ are bounded on $(0, \Delta)$. Since $a(t) \notin L_1(0, \infty)$, (3.21a) gives

$$\phi_1(i\tau) \ge \frac{1}{2} \int_0^{\pi/3\tau} a(t) dt \to \infty, \quad \text{as } \tau \to 0,$$

so that $|\tau/\phi_1(i\tau)|$ is also bounded. Then by (3.26), (3.28), and Corollary 3.2, the

coefficient of $\sin \tau t$ in (3.30) is in $L_1(0, \Delta)$. Using the Riemann-Lebesgue theorem once more, we have

(3.32)
$$\lim_{t \to \infty} y_2(t) = 0.$$

Finally, to treat $y_1(t)$, we note that for $\tau > 0$

$$\frac{d}{d\tau}\phi_1(i\tau) = -\int_0^{\pi/2\tau} ta(t)\sin\tau t\,dt \le 0.$$

Thus by (3.31), $1/\phi_1(i\tau) \downarrow 0$ as $\tau \downarrow 0$; in particular, $1/\phi_1(i\tau)$ is of bounded variation on $[0, \Delta]$, and a familiar theorem concerning the kernel ($\sin \tau t$)/ τ [8, p. 64] yields $y_1(t) \to 0$ as $t \to \infty$. Combining this with (3.29) and (3.32), we have (3.20), and (ii) is proved.

(iii) A proof similar to that for (ii) can be obtained by considering $W_2(s) \equiv W(s) - 1/A(0)s$. We obtain (iii) from (i) as done in the proof of Theorem 2(i) of [5].

By the definition of w(t), we have u(t) = w'(t), where the asymptotic behavior of u(t) is given by conclusion (i). (1.1) for w becomes

(3.33)
$$u(t)-1 = -\int_0^t a(t-\tau)w(\tau) d\tau,$$

since $c + a(t) \in L_1(0, \infty)$ implies c = 0.

When (H3) holds, Levin's proof of Theorem 2(i) in [5] applies word for word to give $w(t) \to (\int_0^\infty a(t) dt)^{-1}$ as $t \to \infty$.

If (H4) holds, we use $A(i\omega) = -i\omega$ (from Lemma 5) and $a(t) \in L_1(0, \infty)$ to compute

$$-\int_0^t a(t-\tau) \frac{2\sin \omega \tau \, d\tau}{3\omega} = \frac{-2}{3\omega} \int_0^t a(\tau) \sin \left[\omega(t-\tau)\right] d\tau$$

$$= \operatorname{Im} \left\{ \frac{-2e^{i\omega t}}{3\omega} \left[A(i\omega) - \int_t^\infty a(\tau)e^{-i\omega \tau} \, d\tau \right] \right\}$$

$$= \frac{2\cos \omega t}{3} + o(1) \qquad (t \to \infty).$$

Then by (3.33) and conclusion (i),

$$\int_0^t a(t-\tau)w_1(\tau) d\tau = 1 + o(1) \qquad (t \to \infty),$$

and $w_1'(t) \to 0$ as $t \to \infty$. The proof of (iii) can now be completed by the method of Theorem 2(i) of [5], which gives $w_1(t) \to (\int_0^\infty a(t) dt)^{-1}$ as $t \to \infty$. This completes the proof of the theorem.

REFERENCES

1. R. Bellman and K. Cooke, Differential-difference equations, Academic Press, New York, 1963.

- 2. G. Doetsch, Theorie und Anwendung der Laplace-Transformation, Springer-Verlag, Berlin, 1937.
 - 3. A. Friedman, On integral equations of Volterra type, J. Analyse Math. 11 (1963), 381-413.
- 4. A. Halanay, On the asymptotic behavior of the solutions of an integro-differential equation, J. Math. Anal. Appl. 10 (1965), 319-324.
- 5. J. J. Levin, The asymptotic behavior of the solution of a Volterra equation, Proc. Amer. Math. Soc. 14 (1963), 534-541.
- 6. J. J. Levin and J. A. Nohel, On a system of integrodifferential equations occurring in reactor dynamics, J. Math. Mech. 9 (1960), 347-368.
- 7. ——, Perturbations of a nonlinear Volterra equation, Michigan Math. J. 12 (1965), 431-447.
 - 8. D. V. Widder, The Laplace transform, Princeton Univ. Press, Princeton, N. J., 1946.

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